

Geometric realisation of Connes spectral triples for algebras with central bases

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Joint work with Prof. Shahn Majid (arXiv:2208.07821)
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WINGs Spring Retreat, Chester eld
18th April 2023

Geometric
realisation of
Connes
spectral
triples for
algebras with
central bases

Preliminaries

NCRG motivation

NCRG formalism

NCRG formalism

Spectral
Triples

Axiomatic
formalism

Comparison with
Connes

Some results for
central bases

Some results for
central bases

Conclusions

1 Preliminaries

- NCRG motivation
- NCRG formalism
- NCRG formalism

2 Spectral Triples

- Axiomatic formalism
- Comparison with Connes
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3 Conclusions

NCRG statement for quantum gravity

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- 1 We will not assume that the spacetime is continuum at the plank scale.
- 2 Instead propose that it is more effectively described by a noncommutative coordinate algebra.
- 3 And in the limit classical RG is recovered.

Basic NCRG formalism

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We work with A a unital algebra, typically a \ast -algebra over \mathbb{C} , in the role of 'coordinate algebra'.

- 1 Differentials are formally introduced as a bimodule¹ of 1-forms equipped with a map $d : A \rightarrow A^1$ obeying the Leibniz rule $d(ab) = (da)b + adb$.
- 2 Assume this extends to an exterior algebra $(\wedge; d)$ with $d^2 = 0$ and d obeying the graded-Leibniz rule and $d(g) = 0$.
- 3 A quantum metric is $g \in \mathcal{L}(A^1, A^1)$ and a bimodule map inverse $(;) : A^1 \rightarrow A^1$.
- 4 A bimodule connection on A^1 is $r : A^1 \rightarrow A^1$ obeying:
 - 1 $r(a \cdot !) = a \cdot r ! + da \cdot !$,
 - 2 $r(! \cdot a) = (r !) \cdot a + (! \cdot da)$
 with a unique 'generalised braiding' bimodule map $\tau : A^1 \otimes A^1 \rightarrow A^1 \otimes A^1$.

In recent years, the quantum Riemannian geometry was extended to a systematic theory including the QLC and further structure

- 1 'spinor' bimodule S equipped with a bimodule connection
- 2 A 'Clifford action': $\sigma^2 = 1$ $\sigma A \sigma^{-1} = S^*$ $\sigma S \sigma^{-1} = S$.
- 3 Leading to a quantum-geometric Dirac operator D .
- 4 And an inner product used to complete to a Hilbert Space.

Motivation for central bases

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We say a basis is central if it is a Grassmann algebra: $e^i e^j = -e^j e^i$.

We will think in central bases if A has trivial centre, S has a central basis f^j and S has a central basis g .

The central bases assumption resumes into the set of 1-forms a_j and a_j adjoint. The metric translates to the matrix metric coefficient in the basis being hermitian. If the ip map then the $-$ -preserving condition translates to the Christoffel symbols in the basis being real.

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Some results and examples

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Theorem

Up to a phase in the 2D spinor bundle case can be obtained with $r > 0$ and $z \in \mathbb{C}$ as either:

$$(1): J = \begin{pmatrix} z & r \\ \frac{1}{r} z^2 & z \end{pmatrix}; \quad (2): J = \begin{pmatrix} 1 & \frac{1}{|z|^2} z \\ z & \frac{z}{z} \end{pmatrix}$$

or its transpose. The $r = 1$ case of (2) needs $z \neq 0$ and up to a

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Noncommutative Torus

the only geometrically realised Dirac operator for the standard Euclidean metric and a WQLC are

$$D(e^i) = (\partial_i + s^i) \cdot e^i + d_i s^i \cdot e^i = i \cdot ((\partial_i + d_i) \cdot e^i) : (1)$$

. With Hilbert space, the state $\sum_{m,n \in \mathbb{Z}} u^m v^n = \sum_{m,n \in \mathbb{Z}} |m\rangle \langle n|$



